# **SELF-DIFFUSION IN MOLECULAR FLUIDS AND NOBLE GASES: AVAILABLE DATA**

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### **ABSTRACT**

Experimental self-diffusivities of gases, vapors, and liquids obtained by means of tracer techniques and nuclear magnetic resonance are reviewed. The considered substances range from noble gases and simple diatomics (nitrogen, oxygen, carbon monoxide, etc.) to complex organic molecules, such as phenolphthalein dimethyl ether and  $2-(\alpha$ -methylbenzylamino)-5nitropyridine, although polymers have not been included. Some comments on the applicability of neutron scattering to the determination of self-diffusion coefficients are also made. All the experimental results of the investigated systems are given as Supporting Information, whereas the references, temperatures, and pressures of these data and the main features of the measurement methods are compiled and classified.

#### **1. INTRODUCTION**

The intradiffusion coefficient can be defined as the translational motion of molecules of a component 1 across a chemically homogeneous mixture of 1 with other substances. The term "self-diffusion"  $(\mathfrak{D}_{11})$  is a special case of intradiffusion where there is only one component in the system.<sup>1</sup> Availability of reliable experimental self-diffusivities allows the testing of theoretical intermolecular potentials for pure substances, since this transport property is easily calculable by computer simulation.<sup>2</sup> Additionally, the development of mathematical models for self-diffusion is interesting from a practical point of view, because they should be capable of predicting binary diffusivities,  $3-6$  which are necessary in the design of industrial processes involving mass transfer.

Unfortunately, reviews of self-diffusion coefficients are scarce in comparison to those of the other two transport properties (viscosity and thermal conductivity), and are included in more general works, dealing with binary (or even multicomponent) diffusivities. In the paper of

Johnson and Babb, $<sup>7</sup>$  that focuses on liquid nonelectrolytes until 1956, only the self-diffusivities</sup> of 11 compounds were cited, whereas in the work of Marrero and Mason,<sup>8</sup> which covered almost all the low-density coefficients prior to 1970, this number increased to 20 (although part of the points were not experimentally obtained but calculated by intermolecular potentials and/or viscosities). The relatively recent book of Winkelmann<sup>9</sup> reported self-diffusivities of 36 fluids (both measured and calculated), including a considerable amount of the data found by Marrero and Mason.

In the present publication, we have attempted to make an overview, as extensive as possible, of all self-diffusion coefficients published to date for liquids, vapors, and gases of organic and inorganic substances, excluding polymers, ionic compounds and metals. More than 15 000 points of 360 substances have been classified in tables according to the measurement techniques, experimental conditions, and other interesting features. Furthermore, all these data are available as Supporting Information, even graphical information. Unlike previous works, we have read directly from the figures with the help of AutoCAD® 2007 when the numeric values were not given in the corresponding articles. The estimated uncertainty for these readings, after testing our procedure in papers where both tables and graphics were present, is around 1-2% (a figure generally lower than experimental precision). On the other hand, self-diffusivities estimated from membrane permeabilities (that change from membrane to membrane<sup>10-12</sup>) or calculated from binary diffusion, intermolecular potentials, or viscosities,13-20 which were cited by Marrero-Mason and/or Winkelmann, have not been considered here.

CAS-numbers, molar masses, melting and boiling points, and critical properties of the compiled substances are also listed in the Supporting Information. They were taken mainly from the handbook of Yaws,<sup>21</sup> but also from the website of the Chemical Abstracts Service<sup>22</sup> and other specific references.23-29

## **2. EXPERIMENTAL PROCEDURES**

**2.1.Tracer methods**. These techniques introduce a small amount of a chemical species which is similar (but not equal) to the molecules of the fluid whose self-diffusion is required, and measure their mobility. The chemical species, called tracer, usually differs from those of the medium by the substitution of one atom by some isotope, i.e., hydrogen by deuterium (D) or tritium (T), carbon 12 by carbon 13 ( ${}^{13}$ C) or 14 ( ${}^{14}$ C), oxygen 16 by oxygen 18 ( ${}^{18}$ O), etc. According to Barton and Speedy,<sup>30</sup> Marrero and Mason,<sup>8</sup> and Winkelmann,<sup>9</sup> the experimental ways to determine the diffusion of the tracer can be classified as closed tube or Loschmidt cell, gas chromatography, gel chromatography, liquid chromatography, interferometry, diffusion bridge, capillary leak, open ended capillary, point source technique, those based on the Dufour effect or on the Kirkendall effect, two-bulb device, back diffusion, diaphragm cell, and the fritted glass plug cell. We have also included three other procedures that were only utilized by their inventors and are not mentioned in those general works, i.e., the sheared boundary apparatus of Devell,<sup>31</sup> the modified capillary method of Pruppacher,<sup>32</sup> and the unsteady-state porous frit cell of Mitchel et al.<sup>33,34</sup>

Tables 1-3 compile all the substances investigated by tracer diffusion, together with the chemical species utilized, the experimental procedure, and the number of data points provided by each reference. When the same reference uses several tracers or methods, this has been divided according to the number or tracers/methods if possible. Self-diffusivities of gases (G), vapors (V), liquids (L), or even supercooled liquids (scL), are commonly given as a function of the temperature  $(T)$  and pressure  $(P)$ , although occasionally, the researchers used the mass density ( $\rho$ ) instead of the pressure. In order to avoid an excessive size of the tables, *P* and  $\rho$  are listed in the same column, and unless specified with the corresponding symbol, numbers have to be interpreted as pressures, because they are the most abundant. Additionally, many liquid selfdiffusivities at atmospheric pressure (101.325 kPa) or at saturation pressure (Sat) only reported the temperature, not the pressure; so, to assign them a condition, we have supposed that diffusion coefficients were all determined at 101.325 kPa if the temperatures are lower than the normal boiling point  $(T_b)$ . Fortunately, this distinction is unimportant, since the values of  $\mathfrak{D}_{11}$  at saturation or at atmospheric pressure in liquids are almost the same below *T*b.

Table 1 is devoted to data given numerically and Table 2 deals with points given in figures only. When the graphical self-diffusivities were correlated with simple equations in *T* and/or *P*, and asterisk (\*) was placed close to the corresponding reference. There are some rare cases in which the data are given exclusively by means of formulas, and these were collected in Table 3. Nonetheless, these mathematical expressions are not of  $\mathfrak{D}_{11}(T,P)$ , but of the intradiffusion coefficients in binary mixtures as a function of the mole fraction (*x*) at a constant *T* and *P*, and therefore, we report only the calculated values for the pure fluid  $(x_1=1.0)$  and not the formulas themselves.

In Table 2, part of the points of Naghizadeh and Rice,  $35$  Annis et al.,  $36$  and Adamson and Irani<sup>37</sup> were those given numerically by Easteal and Woolf,<sup>38</sup> Weissmann and DuBro,<sup>39</sup> and Devell et al., $31$  respectively, but we have not included them in Table 1 since these authors obtained the numbers in the same way as us, i.e., by reading directly from the figures.

When papers do not inform about how many isotopes contained the tracer molecule, only the atom is indicated. In the case of water, we cited tritium and deuterium alone to point out that, although  $D_2O$  or  $T_2O$  are added to  $H_2O$ , the diffusing species are HDO and HTO, respectively.<sup>40</sup> For the krypton, xenon and neon of Groth and co-workers,<sup>41,42</sup> the term "several" has been chosen to designate the tracer, because the authors did not select a single isotope, but separated the gas in a light and a heavy fraction by the Clusius-Dickel technique and caused one fraction to diffuse into the other.

If available, the error (accuracy or precision of the measurements) is also included. Otherwise, the abbreviation n.a. (not available) is written. The very low uncertainties of 0.1% obtained by Vugts and co-workers for argon, nitrogen, carbon monoxide and methane<sup>43-45</sup> and by Trappeniers and Michels for krypton<sup>46</sup> only refer to relative values of  $\mathfrak{D}_{11}(298.16 \text{ K})$  / $\mathfrak{D}_{11}(T)$ or  $\rho\mathfrak{D}_{11}/(\rho\mathfrak{D}_{11})^0$ , respectively, and the error of the absolute self-diffusivities increases to around 0.70%. The superscript "0" indicates that  $\rho$  tends to zero.

On the other hand, most of the procedures cited at the beginning of this section are absolute and do not require calibration. The diaphragm, fritted glass plug, and porous frit cells are relative and always need this, whereas the two-bulb setup can require calibration if the size

of all their parts is not well known, or if it contained a plug of porous material between the chambers. The back diffusion only needs calibration if the exact dimensions are unknown. Liquid self-diffusivities of water,  $40,47$  carbon tetrachloride,  $48$  benzene,  $48,49$  cyclohexane<sup>50</sup> and hexane, or binary diffusivities of carbon tetrachloride in cyclohexane<sup>50</sup> or of heavy water, potassium chloride, sodium chloride and urea in water $51-63$  are the common standards for the three cells, while low density values of carbon dioxide,<sup>64-66</sup> argon,<sup>65,67</sup> hydrogen,<sup>68</sup> carbon dioxide+hydrogen,<sup>69</sup> and hydrogen+deuterium<sup>70</sup> are preferred for the two-bulb device. The back diffusion was only employed once by Harteck and Schmidt<sup>68</sup> for para-hydrogen in normalhydrogen, and the apparatus was calibrated with diffusivities in this mixture obtained by the same authors in a closed tube. Moreover, it has to be said that some researchers did not indicate the source or values of the taken standards,  $7^{1,72}$  that the integral/differential aqueous diffusions of KCl from the several references are all in good agreement except those from the International Critical Tables,  $61,62$  Gordon<sup>59</sup> and Stokes<sup>53,58</sup> beyond 0.3 mol·L<sup>-1</sup> (in fact, Stokes corrected them in a later work<sup>73</sup>) and that Becker et al.<sup>74</sup> did not calibrate their setup and only gave relative results: the absolute values included in the Supporting Information are those obtained years later by Stiel and Thodos,<sup>75</sup> who transformed the points of Becker into absolute ones by comparison with other  $\mathfrak{D}_{11}$  available then.

With regards to the relation between the mobility of the tracer and that of the molecules of the fluid, the kinetic theory establishes that the diffusion of the isotopically labelled species in gases or vapors at low density is<sup>3</sup>

$$
(n\mathfrak{D}_{12})^0 = \frac{3}{8[(1/2)(\sigma_1 + \sigma_2)]^2} \sqrt{\frac{RT(M_1 + M_2)}{2\pi M_1 M_2}} \times \frac{1}{\Omega(\frac{RT}{\sqrt{\varepsilon_1 \varepsilon_2}})}
$$
(1)

where 1 refers to the medium and 2 to the tracer, *n* is the number density,  $\sigma$  is the molecular diameter, *M* is the molar mass, *R* is the gas constant and  $\Omega$  is a function of the temperature and of the characteristic energies  $\varepsilon_1$  and  $\varepsilon_2$ . Assuming that the isotope only changes the mass of the molecule, the self-diffusion can be obtained as

$$
\frac{\mathfrak{D}_{11}}{\mathfrak{D}_{12}} = \sqrt{\frac{2M_2}{M_1 + M_2}}
$$
 (2)

Equation 2 is generally considered valid, and almost all the points at low density in the Supporting Information include this correction.<sup>36,43,64,65,67,76-84</sup> It has been applied even to dense fluids.75,85-89 We have maintained this mass correction in the Supporting Information, except for the methane of Mueller and Cahill<sup>76</sup> and the hydrogen isotopes of Amdur<sup>81</sup> and Mason,  $36,82$  since the former authors converted their data to  $CH_4$  instead to  $CD_4$  and the others transformed the mixtures  $D_2+T_2$  and  $D_2+TH$  to  $H_2$  instead of to  $D_2$ . In the case of liquids, the theoretical knowledge is lower,<sup>90</sup> and although some of the cited references deal with the matter for benzene and cyclohexane<sup>91-96</sup> and a nuclear magnetic resonance study suggested that eq 2 could be applicable too,  $97$  no definitive relation has been accepted and the authors reported the values of  $\mathfrak{D}_{11}$  without correction. A notable exception is that of Mills<sup>40</sup> for water, who measured the

diffusion of HTO in H<sub>2</sub>O, took the values of Longsworth<sup>98,99</sup> for HDO in H<sub>2</sub>O and extrapolated both sets of data to pure  $H_2O$ . He did the same for heavy water as well.

Tracer techniques do not seem to be suitable for liquid helium at very low temperatures, where the quantic effects are so strong that small variations in the amount of  $He<sup>3</sup>$  added to the system modify the measured self-diffusivities greatly.<sup>100-102</sup>



# **Table 1. Self-Diffusion Coefficients of Several Substances, Determined by Tracer Methods, and Reported Numerically**









![](_page_10_Picture_1245.jpeg)

![](_page_11_Picture_873.jpeg)

<sup>a</sup> Original data were only given with the mass correction of eq 2. In the Supporting Information, these corrections have not been applied only in the case of the results of Mueller and Cahill,<sup>76</sup> Amdur and Beatty,<sup>81</sup> and Mason et al.,<sup>36,82</sup> as explained in the text. <sup>*b*</sup>It is not clear if the reported data included the isotopic correction or not. *c*Durbin and Kobayashi<sup>156</sup> employed two devices: the dimensions of the first one were determined by filling it with mercury, but the second one was calibrated with the krypton measurements performed in the first apparatus. <sup>*d*</sup> Points of Wendt et al.<sup>79</sup> were later corrected by Annis et al.,<sup>36</sup> but these latter values were graphically given, and therefore, we have included them in Table 2. <sup>*e*</sup> The critical temperature and pressure for the commercial carbon dioxide utilized were 304.64 K and 74.2 bar, respectively. <sup>*f*</sup>Original values of Mason et al.<sup>82</sup> were later corrected by Annis et al.,<sup>36</sup> and only the corrected points have been included in the Supporting Information. <sup>*g*</sup>Three different apparatuses of the two-bulb type were employed. *h*<sup>1</sup>These measurements of Wang<sup>191,192</sup> were revised and corrected in ref 195.

![](_page_12_Picture_668.jpeg)

**Table 2. Self-Diffusion Coefficients of Several Substances, Determined by Tracer Methods, and Reported Graphically**

<sup>a</sup> The critical temperature of this methane was 190.7 K, the critical pressure was not indicated. <sup>b</sup>The assumed critical constants for argon were 150.87 K and 534.62 kg·m<sup>-3</sup>. <sup>c</sup> Original data were only given with the mass correction of eq 2. <sup>d</sup>These data are those of Wendt et al.<sup>79</sup> following a correction. Two of the nine points are given both numerically and graphically, but for the sake of clarity, all have been included in this table.

| substance   | reference | tracer                  | method     | T/K    | $P/\text{bar}$ | $10^9 \cdot \mathfrak{D}_{11} / (m^2 \cdot s^{-1})$ | state | $%$ error | calibration                                              |
|-------------|-----------|-------------------------|------------|--------|----------------|-----------------------------------------------------|-------|-----------|----------------------------------------------------------|
| benzene     | 113       | ${}^{14}C$              | open ended | 288.80 | 1.01           | 1.963                                               | L     | 1.12      | not required                                             |
| $n$ -hexane | 113       | ${}^{14}C$              | open ended | 288.60 | 1.01           | 3.558                                               | L     | 3.54      | not required                                             |
|             |           |                         |            | 288.80 | 1.01           | 3.710                                               | L     | 2.69      |                                                          |
|             |           |                         |            | 293.50 | 1.01           | 3.932                                               | L     | 2.44      |                                                          |
|             |           |                         |            | 308.20 | 1.01           | 5.082                                               | L     | 1.10      |                                                          |
| cyclohexane | 113       | $^{14}C$                | open ended | 288.60 | 1.01           | 0.926                                               | L     | 4.43      | not required                                             |
|             |           |                         |            | 293.50 | 1.01           | 1.183                                               | L     | 2.28      |                                                          |
|             |           |                         |            | 303.20 | 1.01           | 1.832                                               | L     | 1.47      |                                                          |
| $n$ -hexane | 132       | $C_5{}^{14}CH_{14}$     | diaphragm  | 298.16 | 1.01           | 4.1054 / 4.0955                                     | L     | 0.20      | $0.8 \text{ mol} \cdot \text{L}^{-1}$ urea <sup>52</sup> |
| cyclohexane | 132       | $\rm C_5{}^{14}CH_{12}$ | diaphragm  | 298.16 | 1.01           | 1.4354 / 1.435                                      |       | 0.20      | $0.8 \text{ mol} \cdot \text{L}^{-1}$ urea <sup>52</sup> |

**Table 3. Self-Diffusion Coefficients of Several Substances, Determined by Tracer Methods, and Reported in the Form of Equations**

**2.2.Nuclear Magnetic Resonance**. The estimation of self-diffusion coefficients by nuclear magnetic resonance began with the spin-echo method of  $\text{Hahn}^{224}$  as developed by Carr and Purcell.<sup>225</sup> According to these authors, in the presence of a steady-field gradient (SFG), a 90 $\degree$  pulse at time 0 followed by a single 180 $\degree$  pulse at time  $\tau$  generates a "spin echo" of atomic nuclei at time  $2\tau$ . The amplitudes of the echo immediately after the 90° pulse and at the time  $2\tau$ ,  $A_{(0)}$  and  $A_{(2\tau)}$  respectively, are related by

$$
A_{(2\tau)} = A_{(0)} \exp\left[-\frac{2\tau}{t_{ss}} - \frac{2}{3} \mathfrak{D}_{11} (\gamma G)^2 \tau^3\right]
$$
 (3)

where  $\gamma$  is the gyromagnetic ratio of the studied nucleus,  $t_{ss}$  is the spin-spin relaxation time, and *G* is the magnetic gradient.  $\mathfrak{D}_{11}$  can be derived either from a series of echoes obtained at constant *G* under variation of  $\tau$  or at constant  $\tau$  by applying different gradient values.<sup>226,227</sup>

For measurements of diffusion constants below about  $1 \cdot 10^{-7}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup>, the pulsed field gradient (PFG), proposed by Stejskal and Tanner, is usually employed.<sup>228,229</sup> In this case, *G* is not applied throughout the experiment, but is switched off during the 90º and 180º pulses and during the echo, resulting in

$$
A_{(2\tau)} = A_{(0)} exp\left[-\frac{2\tau}{t_{ss}} - \mathfrak{D}_{11} (\gamma G \delta)^2 \left(\Delta - \frac{\delta}{3}\right)\right]
$$
(4)

being  $\delta$  the duration of the gradient pulse and  $\Delta$  being the time between one pulse and other. When  $\Delta = \delta = \tau$ , the expression of Hahn and Carr-Purcell is recovered.

The main modification of eq 3 is the steady field gradient stimulated echo (SFGSE), which avoids the large loss of the spin-spin relaxation in a viscous fluid during the diffusion time  $\Delta$  by storing the magnetization in the z direction during that time and allows the measurements of  $\mathfrak{D}_{11}$  in supercooled liquids.<sup>230</sup> The pulsed field gradient has been more widely used and has generated numerous improvements, such as the PFG stimulated echo (PFGSE), $^{231}$ the Fourier transform extension of the PFG experiment (PFG-FT),  $232$  the addition of magic angle spinning to the PFG (PFG-MAS),  $^{233}$  the PFG double stimulated echo (PFG-DSTE),  $^{234}$  the alternating PFG  $(A-PFG)$ ,<sup>235</sup> the single-shot method using multiple gradient-recalled echoes  $(SS-MRE)$ ,<sup>236</sup> the spin-echo burst pulse sequence  $(SEBPS)$ ,<sup>237</sup> the diffusion-ordered spectroscopy incorporating bipolar pulse pairs  $(DOSY-BPP)$ ,<sup>238</sup> pulses from a gradient of the radiofrequency field (PRFFG),  $^{239}$  and the one-shot experiment with PRFFG (OS-PRFFG).  $^{240}$  For laser/ thermally polarized <sup>129</sup>Xe at low density, the single-shot PFG stimulated echo sequence  $(SS-PFGSE)^{241}$  and the PFG multiple spin-echo  $(PFGMSE)^{242}$  can be used, whereas for hyperpolarized gaseous  ${}^{3}$ He, the combination of gradient pulses with nuclear magnetic imaging is preferred, both in a simple way (PFG-Im)<sup>243</sup> or in more elaborated sequences (e.g. fast low angle shot, FLASH).<sup>244</sup> Imaging is also applied to self-diffusion of liquids in the form of three different PFG (TDPFG-Im),  $^{245}$  intravoxel incoherent motion imaging (IVIM),  $^{246}$  diffusionweighted contrast-enhanced Fourier acquired steady state technique  $(DW-CE-FAST)$ ,  $247$ 

longitudinal eddy-current delay PFG or longitudinal encode-decode PFG (LED-PFG),<sup>248</sup> quick echo split imaging technique  $(QUEST)$ ,<sup>249</sup> mixed imaging sequence using spin and stimulated echoes (MISSTEC-2 or MISSTEC-4 according to the number of stimulated echoes)<sup>250</sup> and steady-state free precession (SSFP) pulse sequence,<sup>251</sup> although all these latter procedures are mainly focused to medical diagnoses.

Methods of nuclear magnetic resonance that did not use spin echoes have rarely been employed: Jonas and coworkers<sup>252,253</sup> determined the diffusion coefficients in two highly viscous liquids by the procedure of Burnett and Harmon, which is based on the spin-locking field dependence of the rotating-frame proton spin-lattice relaxation time,  $^{254}$  and Barbe et al.  $^{255}$ used the longitudinal relaxation time of  ${}^{3}$ He gas oriented by optical pumping in a nonhomogeneous radiofrequency field.

Tables 4-6 compile all the investigated substances by nuclear magnetic resonance. *P* and  $\rho$  are displayed in one column following the same criteria as Tables 1 and 2, and the asterisks also mean that simple correlations in *T* and/or *P* are available for the graphical data. Some authors distinguish reproducibility from reliability, and then, we have placed the last parameter in italics in brackets. Numerically published alkanes of Bachl,<sup>256</sup> methane of Oosting and Trappeniers,<sup>257</sup> 2-ethylhexylbenzoate of Walker et al.,<sup>253</sup> and ketones and acetates of Petrowsky et al.<sup>258</sup> are published in figures elsewhere,<sup>259-265</sup> so we have only taken into account the numbers and have not added the graphical references to Table 5 nor to the Supporting Information.

If a reference reports diffusivities of the same molecule obtained by different techniques and/or resonant nuclei, we have divided the reference as well. The works of Saitoh et al.<sup>266</sup> and of Asahi and Nakamura<sup>267</sup> were also divided, but for the sake of clarity and not because of procedures or atoms: these researchers measured  $\mathfrak{D}_{11}$  at constant  $\rho$  (in the probe of the nuclear magnetic resonance apparatus) and temperatures above and below the critical one  $(T_c)$ . However, for the subcritical temperatures, most of the conditions correspond to the L-V coexistence zone, the densities are not those initially fixed, and the observed diffusion coefficient is an average value over the liquid and vapor, since the spin echoes from both phases were taken into account (so, we utilized the term "biphasic" to define the state). Notwithstanding this, a part of these coefficients are very close to those of the saturated liquid, due to the fact that the echo of the vapor decays much faster than the liquid except in the vicinity of  $T_c$ .

When there is a peculiarity in the experimental procedure, this has been indicated between parentheses: SS-PFGSE(pol) and FLASH(hyp) point out that the methods have been applied to polarized or hyperpolarized atoms, as previously said; the terms SFG(constant *G*) and  $SFG$ (constant  $\tau$ ) mean that the authors distinguish between points measured at constant *G* under variation of  $\tau$  from those determined at constant  $\tau$  by applying different gradient values; SFG(set 1) and SFG(set 2) take into account the divison of the data in two sets due to accuracy or equipment differences;  $H(-CH_3)$  and  $H(ring)$  refer to the protons of the methyl group of toluene and to the protons of the aromatic ring, respectively; SFG(2.5 mm), SFG(3 mm), SFG(8 mm), SFG(10 mm), SFG(6 cm<sup>2</sup>·g<sup>-1</sup>) and SFG(80 cm<sup>2</sup>·g<sup>-1</sup>) are the internal diameters or specific densities of the probes where Gaines et al.<sup>268</sup> and Hausser et al.<sup>269</sup> placed the samples of helium

3 and water/benzene, respectively, in order to know if the size of these cells restricted the molecular mobility; finally,  $PFG(\Delta=40, G=239)$  and other similar expressions show the time between one pulse and another (in ms) and the magnetic gradient (in gauss/cm) because the authors changed these during the experiments to investigate the influence of both parameters in the self-diffusion of the same compound.

Occasionally, a paramagnetic substance is purposely added to the high-purity chemical in order to shorten the duration of the experiments (by reducing the spin-lattice relaxation time): the xenon of Ehrlich and Carr<sup>270</sup> and of Peereboom et al.<sup>271</sup> contained 0.5% and 0.5% of oxygen, respectively; there was 1.0% O<sub>2</sub> in the ethane of Noble<sup>272,273</sup> and 1.0 g·L<sup>-1</sup> Cu(NO<sub>3</sub>)<sub>2</sub> in the water of Becker et al.<sup>249</sup> Oxygen was also present in the cyclohexane and deuterated aniline of Anderson and Gerritz,<sup>274</sup> whereas the water of Lowe,<sup>235,275</sup> Merbolt<sup>247</sup> and Doran<sup>237</sup> was doped with CuSO4, but none of these latter groups measured or indicated the amount of the corresponding impurity.

Some references included in Table 5 provide figures of  $\mathfrak{D}_{11}$  vs *T* from low temperatures to supercritical conditions, but do not specify the density at which those supercritical selfdiffusivities were obtained. We suspect that it could be the critical one  $(\rho_c)$ , but since we cannot prove this, we have written "Sat and beyond" in the column of the pressure/density.

Table 6 exhibits the self-diffusion coefficients that are only available as mathematical expressions. All of them correspond to liquids and were determined by the SFG, except the carbon tetrachloride of Fischer and Weiss, $^{23}$  that could be measured by PFG as well. We have placed here ten of the twenty-one substances analyzed by Samigullin, $276$  which were also represented graphically, but are not readable due to some deformation in the y-axis of the figures (and detected because a considerable disagreement between the obtained values and those calculated by the formulas). Unlike Table 3, where there were no functions of *T* and/or *P*, all data in this table are given in terms of activation energies  $(E_{\mathfrak{D}})$ , activation volumes  $(V_{\mathfrak{D}})$  and preexponential factors  $(\mathbb{D})$ , defined as

$$
E_{\mathfrak{D}} = -R \left[ \frac{\partial \ln \mathfrak{D}_{11}}{\partial (1/T)} \right]_P \tag{5}
$$

$$
V_{\mathfrak{D}} = -RT \left[ \frac{\partial \ln \mathfrak{D}_{11}}{\partial P} \right]_T \tag{6}
$$

$$
\mathbb{D} = \lim_{\substack{T \to \infty \\ P \to 0}} \mathfrak{D}_{11} \tag{7}
$$

 $\mathbb D$  is commonly replaced by a value of  $\mathfrak{D}_{11}$  at atmospheric pressure and at one reference temperature. If  $E_{\mathcal{D}}$  or  $V_{\mathcal{D}}$  are not constant, we have indicated the temperatures or pressures at which they were determined, so simple interpolations can be done to solve the formulas. In the case of 1,3,5-trimethyl benzene and tetramethyl tin, the activation parameters are complex polynomials in both *T* and *P*, and therefore, in order to spare the reader the integration process, we have put the corresponding  $\mathfrak{D}_{11}(T,P)$  as footnotes.

With regards to the paper of Hertz and Radle,  $277$  who were the first to see the characteristic pressure dependence of water at low temperatures, we have not included this anywhere because results were described in a qualitative way, i.e., they only claimed that the experimental error was around 10%, that  $\mathfrak{D}_{11}$  decreased monotonously with pressure at 298.16 K, and that it went through a maximum at 273.16 K  $(\mathfrak{D}_{11})$  increased from atmospheric pressure to a value (10 to 20) % higher at 147.1 MPa and then decreased to the unspecified atmospheric self-diffusion at 294.2 MPa). Later, other researchers also established the maximum at 298.16 K.

The last column of each table reports the way in which the magnetic gradient is evaluated. The accuracy in the values of  $\mathfrak{D}_{11}$  is dependent on the accuracy in the knowledge of *G*, that is usually estimated in four ways:  $278-282$  (i) theoretically, from the known dimensions, geometry, number of turns of wire in the coil and current applied; (ii) by measuring the changes in the magnetic field when a small sample (commonly, water) is placed at different locations; (iii) by calibration (Cal) with a substance of known self-diffusion; and (iv) by the shape of the spin echo after the 90º pulse, which for cylindrical probes is related to the first-order Bessel function. Dawson et al.<sup>283</sup> and Khoury et al.<sup>284</sup> cited the use of a "flux-gate magnetometer" for this task by a procedure which is very close to the second way; that is, the apparatus indicates the values of the magnetic field along the axis of coils, and *G* is calculated with these data. For the methods of Burnett and Harmon<sup>254</sup> and Barbe et al.,<sup>255</sup> the determination of *G* is not required.

The most employed standard for calibrating  $G$  is the water of Mills,<sup>40</sup> whose selfdiffusivities at (274.16 to 318.16) K and atmospheric/saturation pressure were considered the most reliable and precise ones by Weingärtner in his classical paper of 1982.<sup>285</sup> For higher temperatures and pressures, Weingärtner recommended the data of Krynicky et al.<sup>286</sup> in the ranges (343.2 to 498.2) K, (sat to 170) MPa and of Harris and Woolf<sup>287</sup> at (277.15 to 333.15) K and  $(0.1 \text{ to } 301.8)$  MPa. The tracer diffusivities of benzene of Collings and Mills, <sup>48</sup> available from (288.20 to 318.20) K, are also quite used, but they are preferred for "control", a step which sometimes follows the calibration and consists in measuring the self-diffusion coefficients of other substances whose values are also known in order to ensure the goodness of the calibration step.<sup>259,288-297</sup> Occasionally, the control is carried out with benzene at other temperatures and/or pressures,<sup>298-304</sup> cyclooctane<sup>305,306</sup> and water at different conditions to those of the calibration.304,307 As in Tables 1 and 2, several authors did not indicate the source or values of the taken standards<sup>29,308-314</sup> or the way in which *G* was determined. In the case of Tofts et al.,<sup>315</sup> although the calibration substance is not specified, it is added that the control stage gave values of  $\mathfrak{D}_{11}$  for water very close to those of Mills.

When the researchers that obtained *G* through the shape of the spin echo also reported the H2O self-diffusion (determined in similar form) around 298.16 K, we have added a ratio of that value of water to the self-diffusion of Mills at the same temperature in brackets. If this ratio is far from the unity and is not within the experimental uncertainty, then systematic errors are possible. This is not a general rule, as can be seen for Fischer and Weiss,<sup>23</sup> whose selfdiffusivities for water are very low, but their benzene is in perfect agreement with that of Collings and Mills.<sup>48</sup> The ratio of the water self-diffusivities has been also given when the magnetic gradient was evaluated by the position of a sample.

On the other hand, several question marks can be seen in the three tables. They were placed beside information that was not explicitly given in the references, but reasonably supposed by us. This is the case of Emel'yanov et al.,  $316$  who did not mention the experimental technique, but we attributed the SFG to these authors on the basis that two of them used it in a similar paper published some months later;<sup>317</sup> or of the temperature of Valiev et al.,<sup>318</sup> which was selected because the self-diffusion of water was identical to that measured by them at 293.16 K in another work also of the same year.<sup>319</sup> The temperature at which Van Gelderen et al.<sup>236</sup> performed their experiments was determined after comparison of their points with those of other researchers, and the pressure limit for the alkanes of McCall et al.<sup>320</sup> in Table 6 was established in 55.157 MPa due to the fact that this value, indicated specifically for 2,2 dimethylbutane, was the only value reported throughout the study. In the works of Chien et al.<sup>321</sup> and of Zur et al.<sup>251</sup> we were unable to estimate the temperature at which the experiments were carried out and only the term "room" has been included in the corresponding place in the table.

Additionally, O'Reilly et al.<sup>322</sup> did not mention the value that they obtained for water by the shape of the echo spin in his paper of 1968, but we suppose that it was  $2.23 \cdot 10^{-9}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup> at 298.16 K, since they cited it in subsequent works,<sup>323-325</sup> and is within the 5% of the H<sub>2</sub>O of Simpson and Carr, as the authors said. Oosting and Trappeniers did not mention it either,<sup>257</sup> but the gradient evaluation is so close to that described by them in a previous study,  $326$  where the self-diffusion of water at 298.16 K was  $2.51 \cdot 10^{-9}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup> that we are almost sure that both cases are the same. With regards to the papers of McCall and co-workers, we have assigned the value of 2.33 $\cdot 10^{-9}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup> at 298.66 K<sup>327</sup> to the two studies on paraffins,<sup>320,328</sup> since all the measurements were done with the same equipment within the same year (1958 – 1959). Nonetheless, this is unsuitable for later works, because they affirmed that gradient coils were altered from time to time<sup>329</sup> and reported a value of  $2.45 \cdot 10^{-9}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup> at 298.16 K in 1967.<sup>330</sup> Moreover, it is not clear if the quotation of the diffusivity of Simpson and Carr at room temperature (2.13 $\cdot$ 10<sup>-9</sup> m<sup>2</sup> $\cdot$ s<sup>-1</sup>) in 1959<sup>331</sup> and 1960<sup>332</sup> indicates calibration (except, obviously, for the measurements of water in Table 6) or if they simply selected it thinking that was a more accurate value within their experimental error.

Finally, several misprints and mistakes in the consulted literature have been found: Dupré et al.<sup>333</sup> and Anderson and Gerritz<sup>274</sup> said that they calibrated their devices with the water of Hausser et al.<sup>269</sup> at 293.16 K and with the dodecanol of McCall et al.,<sup>334</sup> respectively, but the points of Hausser are only available above 298.16 K and McCall did not measure  $\mathfrak{D}_{11}$  for dodecanol; the helium diffusivities in the Figure 4 of Luszczynski et al.<sup>335</sup> are multiplied by  $10<sup>2</sup>$ instead of by  $10^3$ , glycerol points of Hrovat and Wade<sup>280</sup> below 313.16 K are also wrongly multiplied by  $10^9$  instead of by  $10^8$ ; some captions of the tables of Bachl<sup>256</sup> are displaced (we identified the correct disposition by comparing the numbers with the figures of the thesis); Panchenkov et al.<sup>336</sup> and Samigullin<sup>276</sup> investigated the dichloroethane, but they did not specify the isomer  $(1,2$ -dichloroethane, 1,1-dichloroethane or a mixture of both); and Fury et al.<sup>337</sup> calibrated their device by erroneously assuming that the pyridine of O'Reilly was pyridine-*d*<sup>5</sup> (in fact, they previously rejected a value of  $1.74 \cdot 10^{-9}$  m<sup>2</sup> $\cdot$ s<sup>-1</sup> at 303.16 K and 0.1 MPa, obtained by evaluating *G* through the shape of the spin echo, which is in good agreement with the

diffusivities of Holz et al.<sup>97</sup> for the deuterated compound), so their results for pyridine- $d_5$  are higher than expected.

| substance      | ref | atom                      | method      | points         | T/K               | P/bar or $\rho$ /(kg·m <sup>-3</sup> ) | state   | % error              | gradient evaluation                                              |
|----------------|-----|---------------------------|-------------|----------------|-------------------|----------------------------------------|---------|----------------------|------------------------------------------------------------------|
| methane        | 338 | H                         | <b>SFG</b>  | 10             | 91.00 to 109.00   | 1.01                                   | L       | 8.10                 | Cal.(water at 298.16 $K^{339}$ )                                 |
|                | 340 | H                         | <b>SFG</b>  | 19             | 298.16 to 364.27  | 67.71 to 173.4                         | G       | 1.00                 | shape                                                            |
|                | 283 | $\boldsymbol{\mathrm{H}}$ | <b>SFG</b>  | 83             | 154.50 to 353.80  | 9.72 to 415.4                          | L, V, G | 6.00                 | flux-gate magnetometer                                           |
|                | 257 | H                         | <b>SFG</b>  | 130            | 90.92 to 307.69   | 9.464 to 452.4 ( $\rho$ )              | L, V, G | $2.00^{341}$         | shape (1.0918? <sup>326</sup> )                                  |
|                | 342 | H                         | <b>SFG</b>  | 65             | 223.16 to 323.16  | 31.50 to 1730.00                       | G       | 2.00                 | Cal.(water <sup>40</sup> and benzene <sup>48</sup> at several T) |
|                | 343 | H                         | <b>SFG</b>  | 36             | 110.00 to 298.15  | 20.76 to 2215.80                       | L, G    | 2.00                 | shape (1.0400)                                                   |
|                | 288 | H                         | PFG-FT      | 45             | 295.0 to 454.00   | 110.00 to 2070.00                      | G       | 1.00 to 2.00 $(5.0)$ | Cal. (water at 298.16 $K^{40}$ )                                 |
|                | 308 | H                         | PFG         | 6              | 303.30, 333.10    | 300.00 to 500.00                       | G       | 5.00                 | Cal.(water)                                                      |
| methane- $d_4$ | 344 | $\mathbf D$               | <b>PFG</b>  | 70             | 143.00 to 454.00  | 210.00 to 2070.00                      | L, G    | 5.00                 | Cal.(heavy water at 298.16 $K^{345}$ )                           |
| ethane         | 288 | H                         | PFG-FT      | 65             | 136.00 to 454.00  | 250.00 to 2000.00                      | L, G    | 1.00 to 2.00 $(5.0)$ | Cal. (water at $298.16 K^{40}$ )                                 |
|                | 308 | H                         | PFG         | 6              | 303.30 and 333.10 | 300.00 to 500.00                       | L, G    | 5.00                 | Cal.(water)                                                      |
| propane        | 288 | H                         | PFG-FT      | 95             | 112.00 to 453.00  | Sat to 2000.00                         | L, G    | 1.00 to 2.00 $(5.0)$ | Cal. (water at 298.16 $K^{40}$ )                                 |
| propane- $d_8$ | 340 | D                         | <b>SFG</b>  | 16             | 298.16 to 358.72  | 42.68 to 173.4                         | L       | 5.00                 | shape                                                            |
| $n$ -butane    | 256 | H                         | <b>SFG</b>  | 49             | 150.00 to 541.00  | Sat to 2000.00                         | L, G    | 3.00(10.0)           | Cal.(water <sup>285</sup> )                                      |
| $n$ -pentane   | 346 | H                         | n.a.        | 51             | 293.00 to 353.00  | 1.00 to 2450.00                        | L       | 4.00 to 5.00         | n.a.                                                             |
|                | 256 | H                         | <b>SFG</b>  | 44             | 145.00 to 450.00  | Sat to 2000.00                         | L       | 3.00(10.0)           | Cal.(water <sup>285</sup> )                                      |
|                | 347 | H                         | PFG-FT      |                | 298.16            | 1.01                                   | L       | n.a.                 | n.a.                                                             |
| $n$ -hexane    | 336 | H                         | <b>SFG</b>  |                | 298.16            | 1.01                                   | L       | 2.43                 | shape (1.0309)                                                   |
|                | 348 | H                         | <b>SFG</b>  |                | 295.16            | 1.01                                   | L       | 5.00                 | n.a.                                                             |
|                | 349 | H                         | <b>SFG</b>  | 59             | 223.15 to 333.15  | 1.00 to 3938.00                        | L       | 1.50(2.5)            | shape (0.9900 to 1.0100)                                         |
|                | 256 | H                         | <b>SFG</b>  | 42             | 188.00 to 455.00  | Sat to 2000.00                         | L       | 3.00(10.0)           | Cal.(water <sup>285</sup> )                                      |
|                | 347 | H                         | PFG-FT      | $\overline{c}$ | 298.16, 323.16    | 1.01                                   | L       | n.a.                 | n.a.                                                             |
|                | 345 | H                         | <b>PFG</b>  |                | 298.16            | 1.01                                   |         | 1.00                 | Cal.(water <sup>40</sup> and benzene <sup>48</sup> at 298.16 K)  |
|                | 308 | H                         | <b>PFG</b>  | 6              | 303.20 and 333.00 | 300.00 to 500.00                       | L       | 5.00                 | Cal.(water)                                                      |
|                | 350 | H                         | <b>SFG</b>  | 10             | 288.16 to 318.16  | Sat to 600.00                          | L       | 3.00                 | Cal.(water + heavy water <sup>207</sup> )                        |
|                | 351 | H                         | <b>PFG</b>  |                | 298.16            | 1.01                                   | L       | 0.40                 | Cal. (water at 298.16 $K^{40}$ )                                 |
|                | 309 | H                         | PFG         | $\overline{4}$ | 298.16 to 313.16  | 1.01                                   | L       | n.a.                 | n.a.                                                             |
|                | 233 | H                         | PFG-MAS     |                | 298.16            | 1.01                                   | L       | n.a.                 | shape (0.9961)                                                   |
|                | 310 | H                         | <b>PFG</b>  | 5              | 303.16 to 323.16  | 1.01                                   | L       | 8.00                 | Cal. (water at several $T$ )                                     |
| $n$ -heptane   | 311 | H                         | <b>PFG</b>  |                | 300.00            | 1.01                                   | L       | n.a.                 | Cal.(water), theory                                              |
|                | 336 | $\, {\rm H}$              | <b>SFG</b>  |                | 298.16            | 1.01                                   | L       | 2.50                 | shape (1.0309)                                                   |
|                | 346 | H                         | n.a.        | 55             | 293.00 to 353.00  | 1.00 to 2450.00                        | L       | 4.00 to 5.00         | n.a.                                                             |
|                | 349 | $\boldsymbol{\mathrm{H}}$ | <b>SFG</b>  | 3              | 195.50 to 299.70  | 1.01                                   | L       | 1.50(2.5)            | shape (0.9900 to 1.0100)                                         |
|                | 256 | H                         | ${\rm SFG}$ | $\,$ 8 $\,$    | 187.00 to 365.00  | Sat                                    | L       | 3.00(10.0)           | Cal.(water <sup>285</sup> )                                      |
|                | 347 | $\boldsymbol{\mathrm{H}}$ | PFG-FT      | $\sqrt{2}$     | 298.16 and 323.16 | 1.01                                   | L       | n.a.                 | n.a.                                                             |
| $n$ -octane    | 336 | $\boldsymbol{\mathrm{H}}$ | <b>SFG</b>  |                | 298.16            | 1.01                                   | L       | 1.45                 | shape (1.0309)                                                   |
|                | 348 | H                         | <b>SFG</b>  |                | 295.16            | 1.01                                   | L       | 5.00                 | n.a.                                                             |

**Table 4. Self-Diffusion Coefficients of Several Substances, Determined by Nuclear Magnetic Resonance, and Reported Numerically**

![](_page_21_Picture_1174.jpeg)

![](_page_22_Picture_1212.jpeg)

![](_page_23_Picture_1083.jpeg)

![](_page_24_Picture_938.jpeg)

![](_page_25_Picture_1080.jpeg)

![](_page_26_Picture_1220.jpeg)

![](_page_27_Picture_1302.jpeg)

![](_page_28_Picture_960.jpeg)

![](_page_29_Picture_941.jpeg)

![](_page_30_Picture_893.jpeg)

![](_page_31_Picture_877.jpeg)

![](_page_32_Picture_1220.jpeg)

![](_page_33_Picture_1265.jpeg)

![](_page_34_Picture_1353.jpeg)

![](_page_35_Picture_681.jpeg)

<sup>*a*</sup> In the Supporting Information, these values, relative to water, have been converted into absolute ones by taking into account the corresponding self-diffusivities of ref 40. <sup>*b*</sup> In the original papers, values of benzene-*d*<sub>6</sub> and pyridine-*d*<sub>5</sub> were multiplied by the square root of the ratio between the molecular mass of the deuterated species and the hydrogenated species, in order to obtain approximated diffusivities for benzene and pyridine. We have undone this correction in the Supporting Information. <sup>*c*</sup>Values of 1,2-dimethoxy ethane, diglyme, and triglyme are higher than others reported graphically by the same authors in subsequent papers.<sup>24,459</sup> According to a private communication of K. Hayamizu,<sup>460</sup> the high diffusivities of this ref 387 are less confidable, since they could be affected by convection. 1,2-diethoxyethane and ethyl propionate are suspected to be erroneous too. <sup>*d*</sup> The nitrobenzene contained 1% H<sub>2</sub>O, since the diffusivity of this small amount of water in the solvent was also studied. <sup>*e*</sup> This chlorobenzene of O'Reilly was erroneous and the correction was given later in the form of an equation,<sup>324</sup> that is cited in Table 6. Tison and Hunt smoothed their experimental measurements taken  $T_c = 318.70$  K and  $\rho_c = 725$  kg·m<sup>-3</sup>. <sup>*g*</sup> Zykov et al.<sup>417</sup> affirm that these data of sulfur hexafluoride are more accurate than others previously published by them,<sup>461</sup> and cited in Table 5. <sup>h</sup> Values of saturated H<sub>2</sub>O and D<sub>2</sub>O of Yoshida et al.<sup>432</sup> were later remeasured and corrected by them in ref 372. <sup>i</sup> This self-diffusion coefficient of Kempka and coworkers is clearly erroneous, since they confirmed it with a point allegedly taken from Kessler et al.<sup>364</sup> which in fact, does not take the value attributed by Kempka.


# **Table 5. Self-Diffusion Coefficients of Several Substances, Determined by Nuclear Magnetic Resonance, and Reported Graphically**





















<sup>*a*</sup> The critical temperature of this ethane doped with 1% O<sub>2</sub> decreased to 305.08 K, but nothing is said about the critical density. <sup>*b*</sup> The critical properties of this ethylene were (282.56±0.05) K,  $(50.66±0.10)$  bar and  $(210±5)$  kg·m<sup>-3</sup>. <sup>c</sup>The critical temperature of this benzene was 563.16 K. <sup>d</sup>The melting point of this 2,2-nitropropane seems appears to be between 326.66 and 327.66 K, a range slightly higher than that found in several databases. <sup>*e*</sup>In the Supporting Information, the relative diffusivities of iodomethane and water of Benedek and Purcell<sup>504</sup> has been converted into absolute ones by taking into account the diffusivities of Mills<sup>40</sup> and Sandhu<sup>501</sup> around 302 K. The same was done with the values of Zykov et al.,<sup>513</sup> which were relatives to the critical diffusivity of sulfur hexafluoride (took as 29.2·10<sup>-9</sup> m<sup>2</sup>·s<sup>-1</sup> from Tison and Hunt,<sup>415</sup> and in good agreement with a subsequent work of Zykov.<sup>417</sup>). *f* O'Reilly said that the critical temperature of the *tert*butyl chloride is 490 K, but all the databases consulted gave a *T<sub>c</sub>* of around 507 K. <sup>*g*</sup> The critical properties of this xenon, doped with 0.5% O<sub>2</sub>, are 289.174±0.003 K and 1097.19 kg·m<sup>-3</sup>. <sup>*h*</sup> The critical properties of this <sup>13</sup>CO<sub>2</sub> were (304.16 ±0.05) K and 468 kg·m<sup>-3</sup>. <sup>*i*</sup> The critical properties for this SF<sub>6</sub> were (318.72±0.05) K, (37.59±0.10) bar and (730±20) kg·m<sup>-3</sup>. *<sup>j</sup>According to* Anderson et al.,<sup>520</sup> the work of Hart has inaccuracies in the temperature scale. <sup>k</sup>Anderson and coworkers measured self-diffusion coefficients of helium 3 three times,<sup>518-520</sup> but considered that in the first case<sup>518</sup> there were systematic errors. <sup>*l*</sup>O'Reilly remeasured the self-diffusivities of hydrogen fluoride and reported them later in the form of an equation,<sup>441</sup> which is compiled in Table 6. *m* Ethanol and propanol of Meckl and Zeidler are given in figures of  $\mathfrak{D}_{11}$  vs  $\rho$ . Fortunately, the paper also provided equations of state that allows us to determine the values of the pressure. *n*The full names of these semi-fluorinated fatty acids were 11,11,12,12,13,13,14,14,15,15,16,16,17,17,18,18,18-heptadecafluoro-oleic acid, 11,1,1,1,2,12,13,13,14,14,15,15,16,16,17,17,18,18,18-18 heptadecafluoro-eleaidic acid and 11,11,12,12,13,13,14,14,15,15,16,16,17,17,18,18,18-heptadecafluoro-stearic acid.



# **Table 6. Self-Diffusion Coefficients of Several Substances, Determined by Nuclear Magnetic Resonance, and Reported in the Form of Mathematical Equations**





<sup>a</sup> The activation volume of ref 542 is a correction of ref 320, so it is expected that the equation will be valid in the same temperature range. <sup>*b*</sup>Also in graphical form. See the text. <sup>*c*</sup> The integrated expression for mesitylene is  $\mathfrak{D}_{11}$ =10<sup>-9</sup>exp(4.994-1343.33/T+2.643/T<sup>2</sup>-0.3867P/T+407.6·10<sup>-6</sup>P+8.2944·10<sup>-8</sup>P<sup>2</sup>-198.4·10<sup>-15</sup>P<sup>3</sup>). <sup>d</sup> The integrated expression for tetramethyl tin is  $\mathfrak{D}_{11}$ =10<sup>-9</sup>exp(4.017-953.37/T+39.444/T<sup>2</sup>-0.0547P/T-590·10<sup>-6</sup>P+14.916·10<sup>-8</sup>P<sup>2</sup>-10.624·10<sup>-15</sup>P<sup>3</sup>). <sup>e</sup> This formula corrects the numerical values published previously by O'Reilly in ref 322. *f* These measurements were considered by O'Reilly better than those given graphically in ref 322. <sup>*g*</sup> Activation energies for diffusion were only provided at 373.16 K. Since the paper apparently demonstrated the applicability of the Stokes-Einstein equation to self-diffusion, we have taken the values at 293.16 K as the energies of the ratio between the viscosity and the absolute temperature, which were reported by the author.

**2.3. Neutron scattering.** When an atomic nucleus of a substance scatters an incident beam of neutrons, there is a response  $S(\vec{Q}, \omega)$  which is a function of the momentum  $\hbar \vec{Q}$  and the energy  $\hbar \omega$  absorbed from or imparted to the neutrons<sup>545,546</sup>

$$
S(\vec{Q}, \omega) = \left\{ \frac{b_{incoh}^2}{b_{total}^2} S_{incoh}(\vec{Q}, \omega) + \frac{b_{coh}^2}{b_{total}^2} S_{coh}(\vec{Q}, \omega) \right\}
$$
(8)

$$
\vec{Q} = \vec{k}_t - \vec{k}_s \tag{9}
$$

$$
\omega = \frac{\hbar}{2m_n} \left( \left| \vec{k_i} \right|^2 - \left| \vec{k_s} \right|^2 \right) \tag{10}
$$

where  $\hbar$  is the rationalized Planck's constant,  $m_n$  the neutron mass,  $\vec{k_i}$  and  $\vec{k_s}$  the wave vector of the incoming and scattered neutrons, respectively (being  $|\vec{k_i}|$  and  $|\vec{k_s}|$  their modules) and *b* the scattering length of the atom. The subscripts "coherent" and "incoherent" refer to the type of scattering,<sup>547</sup> which is coherent if the neutrons are scattered individually by each atom (providing information of their individual motion, i.e., self-diffusion, rotation, vibration, jumps, spin changes, etc.) and is incoherent if there is interference between the scattered waves of the different atoms (showing their collective motion or spatial distribution).

Vanadium and hydrogen are the only naturally occurring elements for which the scattering is predominantly incoherent  $(b<sub>incoh</sub>>b<sub>coh</sub>)$ , but fortunately, the large incoherent scattering length of the protons makes the coherent contribution almost negligible from other nuclei in molecules that contain hydrogen<sup>545</sup> and the scattering spectra will mainly reflect the motions of the individual hydrogen atoms<sup>548-550</sup>

Additionally, the incoherent part can be expressed in terms of the intermediate scattering function  $I_s(\vec{Q}, t)$  as

$$
S_{incoh}(\vec{Q}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} I_s(\vec{Q}, \omega) exp(-i\omega t) dt
$$
 (11)

where translational, rotational, vibrational and all the individual motions are represented. As a working hypothesis, it is commonly supposed that these motions are dynamically independent (or uncoupled) so that the total  $I_s$  can be factorized as the product of separate functions

$$
I_s(\vec{Q}, \omega) = I_s^{trans}(\vec{Q}, \omega) \cdot I_s^{rot}(\vec{Q}, \omega) \cdot I_s^{vibr}(\vec{Q}, \omega) \cdots
$$
 (12)

The main disadvantage of the neutron scattering is that theoretical models have to be assumed for these functions, where not only  $\mathfrak{D}_{11}$  is present, but also a rotational diffusion, characteristics distances, relaxation times and other model-dependent parameters. Moreover, corrections must be made for instrumental resolution, the scattering due to the sample container, multiple scattering and self-shielding $546$ 

The quasielastic part of the spectrum is the peak generated mainly by the translational contribution, and is prominent at low values of the momentum module,  $|\vec{Q}|$ . In the ideal case, the effects of vibrations and rotations are fully negligible, and  $I_s = I_s^{trans}$ . If the atom suffers random alterations of velocity in brief collisions occurring at random times (i.e. obeys Fick's Law), then

$$
I_s^{trans}(\vec{Q}, \omega) = \exp(-\mathfrak{D}_{11}|\vec{Q}|^2 t)
$$
 (13)

and the resulting profile is a simple Lorentzian

$$
S_{incoh}(\vec{Q}, \omega) = \frac{1}{\pi} \frac{\mathfrak{D}_{11}|\vec{Q}|^2}{\omega^2 + (\mathfrak{D}_{11}|\vec{Q}|^2)^2}
$$
(14)

whose width at half-maximum  $(\omega_{1/2})$  verifies that

$$
\omega_{1/2} = 2\mathfrak{D}_{11} |\vec{Q}|^2 \tag{15}
$$

This was the procedure followed by Chen et al.<sup>551</sup> to determine the self-diffusion of supercritical hydrogen at (78 and 293) K and pressures from (13.2 to 197.2) MPa with an accuracy of 10%.

Other researchers, who assumed that eq 15 was valid for liquids (at least, when  $|\vec{Q}|^2 \rightarrow$ 0) and also neglected the non-translational contributions or empirically removed them as background, found  $\mathfrak{D}_{11}$  values close to those derived from tracer methods or nuclear magnetic resonance in the cases of cyclopropane,  $358$  cyclopentene,  $483$  water,  $552,553$  and methanol,  $554$  but observed a considerable disagreement for other compounds, such as pentane,<sup>548</sup> cyclopentane,<sup>550</sup> cyclohexane,<sup>550</sup> toluene,<sup>555</sup> ethanol,<sup>554</sup> pentanol,<sup>554</sup> and glycerol.<sup>548</sup> The reason is that the simple diffusion model is not always obeyed,<sup>556-559</sup> nor even at low values of the momentum module, and that other substance-dependent parameters have to be included in eq 15. Since these parameters change with temperature, and it is not possible to known a priori the applicability of the models and/or simplifications,  $560-562$  the neutron scattering cannot be considered properly a measurement technique, except at certain conditions (e.g., those of Chen and co-workers for hydrogen).

In fact, neutron scattering is a technique that needs the self-diffusion coefficients to validate theories and approximations. The normal procedure when dealing with it is to assign mathematical expressions to all parts of eq 12 and to the coherent contribution, if necessary, and then (i) to obtain  $\mathfrak{D}_{11}$  and other fitting parameters from the experimental spectrum and to compare the calculated self-diffusion coefficients with those available in the literature,<sup>547,562-568</sup> or (ii) to take diffusivities from the literature, to solve eq 8 and to compare the calculated  $S(\vec{Q}, \omega)$  with the experimental one.<sup>569-574</sup> Therefore, in the Supporting Information, we have only included the aforementioned hydrogen of Chen et al.<sup>551</sup>

#### **3. GENERAL CONSIDERATIONS**

In this section, we will make a brief overwiew of the compiled self-diffusivities. It will be mainly focused on general trends and on detecting the strongest discrepancies between the different sources for the same substance. A detailed analysis to determine which experimental points are closer to the real diffusivities is beyond the scope of this study, since not only the data from several researchers can display considerable differences, but also those measured by the same authors in different papers (see, for example the alkanes of von Meerwall,<sup>469,471,473</sup> the alkanes of Emel'yanov,<sup>465,468</sup> the acetone of Le Bihan,<sup>246,389,390</sup> the tri-*n*-butyl phosphate of Pronin and Vashman,  $348,401$ and the 1-octanol and the fatty acids of Iwahashi;<sup>29,347,450,457,458,533,534,541</sup> they, unlike others such as Wang,<sup>191,192,195,198</sup> Hayamizu,<sup>24,387,459,460</sup> Anderson,<sup>518-520</sup> Zykov,<sup>417,461</sup> Yoshida,<sup>372,432</sup> or O'Reilly,<sup>322,441</sup> did not specifically say that some of their measurements were incorrect or inaccurate)

For this comparison, some densities or temperatures of methane, ethylene, saturated xenon, saturated argon, carbon dioxide, and hydrogen which were not provided by the articles themselves are necessary. They were taken from the works of Setzman and Wagner,<sup>575</sup> Smukala et al.,<sup>576</sup> Sifner and Klomfar,<sup>577</sup> Tegeler et al.,<sup>578</sup> Span and Wagner,<sup>579</sup> and Leachman et al.,<sup>580</sup> respectively, changing the critical constants if necessary. These densities or temperatures have been listed in the Supporting Information colored in red. Additionally,  ${}^{13}CO_2$  has been treated as CO<sup>2</sup> because the similarity of their critical properties and the small mass difference (only of  $2.2\%$ ).<sup>27</sup>

**3.1. Water.** As referred to in section 2.2, Weingärtner<sup>285</sup> compared the tracer and spin-echo diffusivities of water available until 1982 and concluded that the extrapolated values of Mills $^{40}$ were the best ones. Previous studies, focused on atmospheric pressure and room temperature<sup>425</sup> or on the pressure dependence at 298  $K^{202}$  had disregarded the points of Wang, <sup>198</sup> Wang et al., <sup>195</sup> Simpson and Carr,  $339$  Trappeniers et al.,  $326$  Benedek and Purcell,  $504$  Cuddeback et al.,  $194$ Kisel'nik et al.,<sup>526</sup> and McCall et al.<sup>331</sup> Subsequent works, such as those of Harris and Newitt<sup>430</sup> and Yoshida et al.<sup>372</sup> confirmed the good agreement between their measurements and the data of Angell et al.,<sup>527</sup> Harris and Woolf,<sup>287</sup> Hausser et al.,<sup>269</sup> Prielmeier et al.,<sup>293</sup> and Krynicki et al.,<sup>286</sup> although the pressurized diffusivities of Krynicki below 298.16 K are erroneous. The temperature dependence of  $\mathfrak{D}_{11}$  at 101.325 kPa is also correct for Ertl and Dullien,<sup>378</sup> Becker et al.,<sup>249</sup> Holz et al.,<sup>307</sup> Tofts et al.,<sup>315</sup> and Hardy et al.<sup>529</sup>

Therefore, in those papers of the following sections that calibrated the coils with the water of Simpson and Carr  $(2.13 \cdot 10^{-9} \text{ m}^2 \cdot \text{s}^{-1})$  at 298.16 K) or of Trappeniers et al.  $(2.51 \cdot 10^{-9} \text{ m}^2 \cdot \text{s}^{-1})$  $m^2 \cdot s^{-1}$  at 298.16 K), we have divided the corresponding  $\mathfrak{D}_{11}$  values by 0.9265 or 1.0918, respectively, in order to obtain self-diffusion coefficients as if the apparatuses had been calibrated with the water of Mills  $(2.299 \cdot 10^{-9} \text{ m}^2 \cdot \text{s}^{-1})$  at 298.16 K). These corrections have not been included in the Supporting Information, and have not been applied to those studies which

evaluated *G* by the shape of the echo-spin either, since in this case, the presence of systematic error is not always assured (see Section 2.2).

With regards to heavy water, it has to be said that, within the experimental uncertainties, all the liquid diffusivities at atmospheric/saturation pressure are in relatively good agreement, except those of Wilbur et al.<sup>435</sup> around (363.16 to 393.16) K (the uncorrected ones of Yoshida et al.<sup>432</sup> have not been taken into account). At high pressures, the concordance gets worse, especially between Prielmeier-Arnold<sup>293,438</sup> and DeFries and Jonas<sup>436</sup> in the range (258.16 to 278.16) K and between Wilbur et al.<sup>435</sup> and Prielmeier et al.<sup>293</sup> at (363.16 to 423.16) K.

**3.2. Other saturated liquids.** For these comparisons, extrapolation of some highpressure diffusivities to low pressure has been carried out.<sup>35,108,412,413</sup>

At first glance, all the coefficients of Kitchlew and Nageswara  $Rao<sup>484</sup>$  are systematically higher than those determined by other authors for the same substances (assuming that the equations of Samigullin for *m*-xylene and *p*-xylene are valid in the same temperature ranges). The values obtained for hexane by Vashman and Pronin,<sup>348</sup> for dodecane by Emel'yanov et al.,<sup>468</sup> for tetratriacontane by Yamakawa et al.,<sup>354</sup> for dioxane by Clemett,<sup>491</sup> for dimethylsulfoxide by Merbolt et al.,  $^{247}$  for tri-*n*-butyl phosphate by Pronin and Vashman,  $^{401}$  for fluoroform by Chaffin and Hubbard,<sup>403</sup> for *tert*-butyl chloride by Grochulski et al.,<sup>496</sup> for xenon by Yen and Norberg,<sup>420</sup> for carbon dioxide by Krynicki,<sup>512</sup> for methanol by Sandhu,<sup>501</sup> and for ethanol, propanol and butanol by Powles and Cutler<sup>443</sup> are very high when also compared with the measurements from other researchers. On the other hand, the methane of Gaven et al., $^{226}$  the decane of Emel'yanov et al.,  $468$  the toluene of Reimschüssel and Hawlicka,  $217$  the acetone determined by TDPFG-Im<sup>245</sup> and SS-MRE,<sup>236</sup> and the xenon of Streever and Carr<sup>410</sup> are considerable lower than the rest.

It can also be observed that for hexane, octane, decane, tetradecane, tetracosane, neopentane, isopentane, 2,2-dimethylbutane, nitromethane, *N,N*-dimethylformamide, dimethylsulfoxide, dichloromethane, chloroform, *tert*-butyl chloride, hexafluorobenzene, chlorobenzene, octanol, stearic acid, oleic acid, and eleaidic acid, there was good concordance between the several sources at low temperatures, but at higher ones, some sets of diffusivities display higher activation energy than the others. See, for example, Figure 1. Near the critical temperature, the cyclohexane, benzene and toluene of Panchenkov et al.<sup>365</sup> are (20 to 30) % higher than those of Yoshida et al.,<sup>372</sup> Hausser et al.,<sup>269</sup> Asahi et al.,<sup>267</sup> and Shimokawa,<sup>486</sup> although at lower temperatures, the data of Yoshida, Hausser and Asahi are 10% above the values of Panchenkov. These discrepancies are frequently attributed to convection (especially when nuclear magnetic resonance is utilized) which causes measured diffusivities to be above the real ones,  $351,364,489,581,582$  but this fact does not guarantee that the lower values are always the correct ones: some investigators who perceived that their points were higher than those expected did not detect convection in their devices $420,268$  and other experimental errors can decrease the experimental results below the true self-diffusion coefficients.<sup>125,372,432</sup>



**Figure 1.** Self-diffusion of hexane at saturation/atmospheric pressure as a function of the temperature: open rectangle, Hawlicka and Reimschüssel;<sup>113</sup> A, Bachl;<sup>256</sup> O, Brüsewitz and Weiss;<sup>301,302</sup>  $\bullet$ , Kasahara et al.;<sup>310</sup> +. Douglass and McCall;<sup>328</sup> **.** Iwahashi et al.;<sup>347</sup>  $\blacklozenge$ . Vashman and Pronin;<sup>348</sup>  $\triangle$ . Harris;<sup>349</sup>  $\times$ . Marbach and Hertz; $^{350}$   $\star$ , Agishev and Emel'yanov; $^{465}$   $\diamond$ , Panchenkov et al.; $^{467}$   $\Box$ , Emel'yanov et al. $^{468}$ 

For cyclohexanone, tetramethyl lead, bromoform, iodomethane, and krypton the differences among the few available diffusivities is so strong, that it is not possible to estimate  $\mathfrak{D}_{11}$  with less than 20% error even at a single temperature.

Based on the experimental arrangements, Mills<sup>425</sup> concluded that, at 298.16 K, the argon of Cini Castagnoli<sup>154</sup> was more accurate than that of Corbett and Wang<sup>153</sup> and of Naghizadeh and Rice.<sup>35</sup> By comparing calculated activation energies, Freer and Sherwood<sup>125</sup> stated that the cyclohexane data of McCall et al.,<sup>331</sup> Hawlicka and Reimschüssel,<sup>113</sup> and Anderson and Gerritz<sup>274</sup> were less reliable than those of others, and said the same for the benzene of Graupner and Winter,  $^{126}$  Hiraoka et al.,  $^{128}$  and McCall et al.<sup>331</sup>

**3.3. Compressed liquids**. Ethane of Wade and Waugh,<sup>463</sup> benzene of Hiraoka et al.<sup>128</sup> and McCall et al., $331$  carbon tetrachloride of Watts et al., $145$  carbon disulfide of Koeller and

Drickamer,<sup>181</sup> and methanol of Hiraoka et al.<sup>212</sup> have been critizised and considered erroneous by Greiner-Schmid et al.,<sup>288</sup> McCool et al., <sup>133</sup> McCool and Woolf,<sup>148</sup> Woolf,<sup>72</sup> and Hurle et al.,<sup>583</sup> respectively. The concordance among the rest of sources for each liquid is relatively good, except for the alkanes of McCall and coworkers,<sup>320,328,331,542</sup> the methanol-*d* of Jonas and Akai above 283 K, $446$  the 2,2-dimethylbutane, and the cyclohexane. In the first case, the agreement with the paraffins of Harris,  $137,349$  Bachl,  $256$  Arkhipov et al.,  $346$  Marbach and Hertz,  $350$  and Brüsewitz and Weiss $302$  is only good at certain temperatures, and for the isopentane of Enninghorst<sup>356</sup> there is no agreement at all. In the case of methanol- $d$ , the high-temperature points of Jonas and Akai are (12 to 25) % higher than those of Hurle et al.<sup>213</sup> and of Karger et al.<sup>295</sup> For 2,2-dimethylbutane and cyclohexane, each author obtained a different trend, as shown in Figures 2 and 3, respectively.



**Figure 2.** Temperature and pressured dependence of self-diffusion of 2,2-dimethyl butane according to three sources: O, Bachl;<sup>256</sup>  $\times$ , Polzin and Weiss;<sup>304</sup> solid line, McCall et al.<sup>320</sup> For the sake of comparison, some points have been interpolated.



**Figure 3.** Temperature and pressure dependence of the self-diffusion of cyclohexane according to four sources:  $\Box$ , McCool and Woolf;<sup>124</sup> ×, Polzin and Weiss;<sup>304</sup> solid line, McCall et al.;<sup>331</sup> A, Jonas et al.<sup>367</sup> For the sake of comparison, some points have been interpolated.

On the other hand, and within an experimental error of 10%, there is no difference between the self-diffusivity of normal hydrogen and a mixture with other proportions of *ortho*and *para*-hydrogen at constant temperature and density.<sup>515</sup>

**3.4. Vapors and gases at low density.** For each substance, all the points from the different sources are in relatively good agreement (5 to 10%), except the value of Hutchinson for argon at 90.16  $K^{67}$  and the xenon of Benenson et al.<sup>163</sup> and of Groth and Harteck.<sup>41</sup> Nonetheles, Hutchinson and Benenson themselves doubt the reliability of these measurements.

Mueller and Cahill<sup>76</sup> questioned their acetylene also because experimental problems. Unfortunately, no more data from other investigators are available in the same conditions to confirm or deny their suspicions.

Additionally, it has to be pointed out that the high-temperature data of Ember et al.<sup>175</sup> and Pakurar and Ferron<sup>177</sup> were cast in doubt by Bock et al.,<sup>584</sup> who were able to reproduce the low-density viscosity of carbon dioxide within a 1% error for (1000 to 1500) K by calculations with the intermolecular potential of Bukowski, but obtained deviations of (13 to 43) % for the self-diffusivities of the aforementioned researchers.

The applicability of eq 2 to the hydrogen isotopes has been discussed by Amdur and Beatty, $81$  Mason et al.  $36,82$  and Reichenbacher et al.  $187$  The first two groups observed differences of (4 to 6) % and stated that not only the mass, but also the characteristic energies and molecular diameters changed considerably; whereas the last group affirmed that, although eq 2 was not strictly valid, the deviations were around 2% and attributed the results of Amdur and Mason to experimental errors.

As in the case of the compressed liquid,  $\rho \mathfrak{D}_{11}$  for hydrogen is independent of the concentration of the ortho species at a given temperature.<sup>515</sup>

**3.5. Compressed gases and vapors.** Previous studies<sup>27,585</sup> have rejected the methane of Cini Castagnoli et al.<sup>216</sup> and the carbon dioxide of Duffield and Harris<sup>89</sup> since these authors were the only ones to see an anomaly of the self-diffusion (similar to that of viscosity and thermal conductivity) in the vicinity of the critical point. Other researchers<sup>412,586</sup> have considered erroneous the methane and carbon dioxide of Drickamer and coworkers $86,104,173,174$  as well, but because their points were 50% lower than the rest of the references.

In order to afford a better comparison among the literature, we have plotted  $\rho \mathfrak{D}_{11}$  vs  $\rho$ for some substances, as in Figure 4. At first glance, it is clear that the ethylene and sulfur hexafluoride of Hamann<sup>476,477</sup> and the carbon dioxide of Krynicki et al.<sup>512</sup> are 20-40% higher than those from other sources, although if the values of Hamman are divided by 1.1091 (see Table 5), the ethylene becomes closer to the data of Takahashi,  $^{118}$  Arends,  $^{357}$  and Peereboom.<sup>479</sup> This improvement can be also appreciated in Figure 4, where the data of Oosting and Trappeniers<sup>257</sup> around 298 K, divided by 1.0918, approach the more recent values of Harris<sup>342</sup> and Greinerd-Schmidt.<sup>288</sup> Nonetheless, the quality of the correction is not constant, since at low temperatures, the concordance between Harris<sup>343</sup> and Oosting<sup>257</sup> is good without taking into account this factor.

On the other hand, it is curious that at 298 K, all the diffusivities obtained by nuclear magnetic resonance tend to a low-density value below that of Winn,  $65$  which is in accordance with the extrapolation of Takahashi<sup>107</sup> and the measurements in methane- $d_4$  of Mueller<sup>76</sup> and Vugts.<sup>45</sup> However, this cannot be due to differences between the tracer motion and the spin echo, since the tracer chromatographic method of Hu and Kobayashi<sup>106</sup> provides the same result as the nuclear magnetic resonance, and at lower temperatures, the agreement among the several

data sets is better, all them merging into the single point of Winn around 194 K. We have seen the same discrepancies regarding the low-density limit for the carbon dioxide of Etesse et al. $412$ and for the hydrogen of Chen et al.<sup>419</sup> as well, but not for ethylene or xenon.



**Figure 4.** Density dependence of the group  $\rho \mathfrak{D}_{11}$  for methane according to several sources:  $\bullet$ , Winn;<sup>65</sup>  $\blacklozenge$ , Hu and Kobayashi;<sup>106</sup> +, Takahashi;<sup>107</sup>  $\Box$ , Dawson et al.;<sup>283</sup>  $\times$ , Oosting and Trappeniers;<sup>257</sup>  $\odot$ , Greiner-Schmid et al.;<sup>288</sup>  $\triangle$ , Harris.<sup>342,343</sup> For the sake of comparison, some points have been interpolated. The dotted line corresponds to the data of Oosting and Trappeniers divided by the factor 1.0918.

In the case of carbon tetrafluoride, where only two sets of supercritical points are available, it is not possible to know which are correct, since those of Khoury and Kobayashi<sup>284</sup> are systematically (20 to 40) % lower than those of Has and Lüdemann.<sup>291</sup> For hydrogen, the neutron scattering<sup>551</sup> produced results (5 to 10) % higher than the spin-echo technique,<sup>419</sup> the former being in better agreement with the low-density tracer diffusivities than the latter.

### **4. CONCLUSIONS**

Self-diffusion in fluids has been reviewed. The most widely studied substance is water, followed by benzene and cyclohexane. The preferred measurement technique is the spin-echo, in the pulsed field gradient version of Stejskal-Tanner or in one of its numerous improvements, since it does not need the introduction of a new species in the system (tracer methods) nor the assumption of complex mathematical models for the molecular motions (neutron scattering). Additionally, we hope that the thousands of experimental data compiled as Supporting Information will facilitate the understanding of this fundamental transport property for future researchers.

### **ASSOCIATED CONTENT**

#### **Supporting information**

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jced.5b00323.

Properties of the considered substances and all the data of Tables 1,2,4, and 5, together with the points obtained by Chen et al.551 for hydrogen through neutrón scattering (ZIP)

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## **Notes**

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## **SELF-DIFFUSION IN MOLECULAR FLUIDS AND NOBLE GASES: AVAILABLE DATA**

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